

## Solutions to Linear analysis, February 01, 2017

1. The cosine series coefficients are

$$a_0 = \frac{2}{\pi} \int_0^{2\pi} f(x) dx = \frac{2}{\pi} \int_{\pi-a}^{\pi} dx = \frac{2a}{\pi}$$

and

$$\begin{aligned} a_n &= \frac{2}{\pi} \int_0^{\pi} f(x) \cos nx dx = \frac{2}{\pi} \int_{\pi-a}^{\pi} \cos nx dx = \frac{2}{\pi} \left[ \frac{\sin nx}{n} \right]_{\pi-a}^{\pi} = \\ &= \frac{2}{\pi n} (\sin n\pi - \sin(n\pi - na)) = \frac{2}{\pi} (-1)^n \frac{\sin na}{n}, \quad n \geq 1. \end{aligned}$$

Hence

$$(1) \quad f(x) \sim \frac{a}{\pi} + \frac{2}{\pi} \sum_{n=1}^{\infty} (-1)^n \frac{\sin na}{n} \cos nx$$

is the cosine series expansion of  $f(x)$  on  $0 \leq x \leq \pi$ . Let  $g(x)$  be the  $2\pi$ -periodic function on  $\mathbf{R}$  for which  $g(x) = 0$  for  $0 \leq x \leq \pi - a$ ,  $g(x) = 1$  for  $\pi - a < x < \pi - a$  and  $g(x) = 0$  for  $\pi + a \leq x < 2\pi$ . Then  $g(x) = f(x)$  for  $0 \leq x \leq \pi$  and  $g(x)$  is an even function on  $\mathbf{R}$ . It follows that the Fourier series expansion of  $g(x)$  on  $\mathbf{R}$  is  $g(x) \sim \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos nx$ ,  $x \in \mathbf{R}$ , with the same coefficients  $a_n$ ,  $n = 0, 1, \dots$  as above. The function  $g(x)$  is piecewise  $C^1$  on  $\mathbf{R}$ . The Dirichlet convergence theorem for Fourier series therefore shows that the Fourier series of  $g(x)$  converges for all  $x \in \mathbf{R}$  and that its sum is equal to  $\frac{1}{2}(g(x+) + g(x-))$  for all  $x \in \mathbf{R}$ . That is  $\sum_{n=1}^{\infty} a_n \cos nx = \frac{1}{2}(g(x+) + g(x-)) - \frac{1}{2}a_0$  for all  $x \in \mathbf{R}$ . In particular

$$(2) \quad \sum_{n=1}^{\infty} a_n \cos nx = \begin{cases} -\frac{a}{\pi} & \text{if } 0 \leq x < \pi - a \\ \frac{1}{2} - \frac{a}{\pi} & \text{if } x = \pi - a \\ 1 - \frac{a}{\pi} & \text{if } \pi - a < x < \pi + a \\ \frac{1}{2} - \frac{a}{\pi} & \text{if } x = \pi + a \\ -\frac{a}{\pi} & \text{if } \pi + a < x < 2\pi. \end{cases}$$

We now make repeated use of a theorem of integration of Fourier series expansion. The expansion  $g(x) \sim \frac{1}{2}a_0 + \sum_{n=1}^{\infty} a_n \cos nx$ ,  $x \in \mathbf{R}$ , implies

$$\int_0^x g(t) dt = \int_0^x \frac{1}{2}a_0 dt + \sum_{n=1}^{\infty} a_n \int_0^x \cos nt dt \quad \text{for all } x \in \mathbf{R} \quad \iff$$

$$\int_0^x g(t) dt = \frac{a}{\pi}x + \sum_{n=1}^{\infty} \frac{a_n}{n} \sin nx \quad \text{for all } x \in \mathbf{R}.$$

Set  $h(x) = \int_0^x g(t) dt$ ,  $x \in \mathbf{R}$ . Then  $h(x) - \frac{a}{\pi}x$  is  $2\pi$ -periodic since  $h(x) - \frac{a}{\pi}x = \sum_{n=1}^{\infty} \frac{a_n}{n} \sin nx$  for all  $x \in \mathbf{R}$ . Calculation gives

$$h(x) = \begin{cases} 0 & \text{if } 0 \leq x \leq \pi - a \\ x - (\pi - a) & \text{if } \pi - a < x < \pi + a \\ 2a & \text{if } \pi + a \leq x < 2\pi. \end{cases}$$

The expansion  $h(x) - \frac{a}{\pi}x = \sum_{n=1}^{\infty} \frac{a_n}{n} \sin nx$ ,  $x \in \mathbf{R}$ , implies

$$\int_0^x h(t) dt - \int_0^x \frac{a}{\pi}t dt = \sum_{n=1}^{\infty} \frac{a_n}{n} \int_0^x \sin nt dt \quad \text{for all } x \in \mathbf{R} \quad \Longleftrightarrow$$

$$\int_0^x h(t) dt - \frac{a}{2\pi}x^2 = b - \sum_{n=1}^{\infty} \frac{a_n}{n^2} \cos nx \quad \text{for all } x \in \mathbf{R}$$

where  $b = \sum_{n=1}^{\infty} \frac{a_n}{n^2}$ . Set  $k(x) = \int_0^x h(t) dt$ ,  $x \in \mathbf{R}$ . Then  $k(x) - \frac{a}{2\pi}x^2$  is  $2\pi$ -periodic since  $k(x) - \frac{a}{2\pi}x^2 = b - \sum_{n=1}^{\infty} \frac{a_n}{n^2} \cos nx$  for all  $x \in \mathbf{R}$ . Calculation gives

$$(3) \quad k(x) = \begin{cases} 0 & \text{if } 0 \leq x \leq \pi - a \\ \frac{1}{2}(x - (\pi - a))^2 & \text{if } \pi - a < x < \pi + a \\ 2a(x - \pi) & \text{if } \pi + a \leq x < 2\pi. \end{cases}$$

Finally integrating the expansion  $k(x) - \frac{a}{2\pi}x^2 = b - \sum_{n=1}^{\infty} \frac{a_n}{n^2} \cos nx$ ,  $x \in \mathbf{R}$ , over the interval  $0 \leq x \leq 2\pi$  we get

$$\int_0^{2\pi} k(x) dx - \int_0^{2\pi} \frac{a}{2\pi}t^2 dt = \int_0^{2\pi} b dx - \sum_{n=1}^{\infty} \frac{a_n}{n^2} \int_0^{2\pi} \cos nx dx \quad \Longleftrightarrow \quad b = \frac{1}{2\pi} \int_0^{2\pi} k(x) dx - \frac{2\pi a}{3}.$$

Calculation gives  $b = -\frac{1}{6\pi}a(\pi^2 - a^2)$ . Consequently

$$\sum_{n=1}^{\infty} \frac{a_n}{n^2} \cos nx = \frac{a}{2\pi}x^2 - \frac{1}{6\pi}a(\pi^2 - a^2) - k(x), \quad x \in \mathbf{R}.$$

Together with (3) this gives the sum of  $\sum_{n=1}^{\infty} \frac{a_n}{n^2} \cos nx$  for  $0 \leq x < 2\pi$  (and therefore for all  $x \in \mathbf{R}$  since the sum is  $2\pi$ -periodic).

Using (2) for  $x = \pi$  gives  $\sum_{n=1}^{\infty} (-1)^n a_n = 1 - \frac{a}{\pi} \Longleftrightarrow \sum_{n=1}^{\infty} \frac{\sin na}{n} = \frac{1}{2}(\pi - a)$ . Here  $0 < a < \pi$ . Clearly  $\sum_{n=1}^{\infty} \frac{\sin nu}{n} = 0$  for  $u = 0$ . We also note that  $\sum_{n=1}^{\infty} \frac{\sin nu}{n}$  is an odd  $2\pi$ -periodic function. Let  $\alpha(u)$  be the  $2\pi$ -periodic function for which  $\alpha(0) = 0$  and  $\alpha(u) = \frac{1}{2}(\pi - u)$  for  $0 < u < 2\pi$ . Then  $\alpha(u)$  is an odd  $2\pi$ -periodic function which is equal to  $\sum_{n=1}^{\infty} \frac{\sin nu}{n}$  for  $0 \leq u < \pi$ . It follows that  $\sum_{n=1}^{\infty} \frac{\sin nu}{n} = \alpha(u)$  for all  $u \in \mathbf{R}$ . In particular  $\sum_{n=1}^{\infty} \frac{\sin nu}{n} = 0$  for  $u = 0$  and  $\sum_{n=1}^{\infty} \frac{\sin nu}{n} = \frac{1}{2}(\pi - u)$  for  $0 < u < 2\pi$ .

Applying the Parseval identity to the expansion (1) gives

$$\frac{1}{2} \left( \frac{2a}{\pi} \right)^2 + \frac{4}{\pi^2} \sum_{n=1}^{\infty} \frac{\sin^2 na}{n^2} = \frac{2}{\pi} \int_0^{\pi} |f(x)|^2 dx.$$

But  $\int_0^{\pi} |f(x)|^2 dx = a$ . Thus  $\sum_{n=1}^{\infty} \frac{\sin^2 na}{n^2} = \frac{1}{2}a(\pi - a)$ . Here  $0 < a < \pi$ . Clearly  $\sum_{n=1}^{\infty} \frac{\sin^2 nu}{n^2} = 0$  for  $u = 0$ . We also note that  $\sum_{n=1}^{\infty} \frac{\sin^2 nu}{n^2}$  is a  $\pi$ -periodic function since  $\sin^2 n(\pi + u) = \sin^2 nu$  for all  $u \in \mathbf{R}$ . Let  $\beta(u)$  be the  $\pi$ -periodic function for which  $\beta(u) = \frac{1}{2}u(\pi - u)$  for  $0 \leq u < \pi$ . It follows then that  $\sum_{n=1}^{\infty} \frac{\sin^2 nu}{n^2} = \beta(u)$  for all  $u \in \mathbf{R}$ . In particular  $\sum_{n=1}^{\infty} \frac{\sin^2 nu}{n^2} = \frac{1}{2}u(\pi - u)$  for  $0 \leq u < \pi$  and  $\sum_{n=1}^{\infty} \frac{\sin^2 nu}{n^2} = \frac{1}{2}(u - \pi)(2\pi - u)$  for  $\pi \leq u < 2\pi$ .

2. The Fourier transform of  $g(x)$  is

$$\begin{aligned}\hat{g}(u) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} g(x) e^{-iux} dx = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-|x|} e^{-iux} dx = \\ &= \frac{1}{2\pi} \int_0^{\infty} e^{-x} (e^{-iux} + e^{iux}) dx = \frac{1}{2\pi} \int_0^{\infty} (e^{-(1+iu)x} + e^{-(1-iu)x}) dx = \\ &= \frac{1}{2\pi} \left[ \frac{e^{-(1+iu)x}}{-(1+iu)} + \frac{e^{-(1-iu)x}}{-(1-iu)} \right]_0^{\infty} = \frac{1}{2\pi} \left( \frac{1}{1+iu} + \frac{1}{1-iu} \right) = \frac{1}{\pi(1+u^2)}.\end{aligned}$$

The function  $g(x)$  is continuous on  $\mathbf{R}$  and has a piecewise continuous derivative on  $\mathbf{R}$ . Also  $\int_{-\infty}^{\infty} |\hat{g}(u)| du < \infty$ . Thus

$$\int_{-\infty}^{\infty} \frac{1}{\pi} \frac{1}{1+u^2} e^{ixu} du = e^{-|x|} \quad \text{for all } x \in \mathbf{R}$$

by the Fourier inversion formula. Equivalently

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{1+u^2} e^{-ixu} du = \frac{1}{2} e^{-|x|} \quad \text{for all } x \in \mathbf{R}.$$

Consequently  $f(x)$  has the Fourier transform  $\hat{f}(u) = \frac{1}{2} e^{-|u|}$ . The function sequence  $f_n(x)$   $n = 1, 2, \dots$  on  $\mathbf{R}$  is defined by:  $f_1(x) = f(x)$  and  $f_k(x) = (f * f_{k-1})(x)$  for  $k = 2, 3, \dots$ . Applying the Fourier transform to the formula  $f_k(x) = (f * f_{k-1})(x)$ ,  $k \geq 2$ , gives  $\hat{f}_k(u) = 2\pi \hat{f}(u) \hat{f}_{k-1}(u)$ ,  $k \geq 2$ . Now let the integer  $n \geq 1$  be arbitrary. Repeated application of the formula  $\hat{f}_k(u) = 2\pi \hat{f}(u) \hat{f}_{k-1}(u)$ ,  $k \geq 2$ , gives

$$\begin{aligned}\hat{f}_n(u) &= 2\pi \hat{f}(u) \hat{f}_{n-1}(u) = 2\pi \hat{f}(u) 2\pi \hat{f}(u) \hat{f}_{n-2}(u) = (2\pi)^2 (\hat{f}(u))^2 \hat{f}_{n-2}(u) = \dots \\ &= (2\pi)^{n-1} (\hat{f}(u))^{n-1} \hat{f}_1(u) = (2\pi)^{n-1} (\hat{f}(u))^{n-1} \hat{f}(u) = (2\pi)^{n-1} (\hat{f}(u))^n.\end{aligned}$$

Hence

$$\hat{f}_n(u) = (2\pi)^{n-1} (\hat{f}(u))^n = (2\pi)^{n-1} \left( \frac{1}{2} e^{-|u|} \right)^n = \pi^{n-1} \frac{1}{2} e^{-|nu|} = \pi^{n-1} \hat{f}(nu)$$

and therefore

$$f_n(x) = \pi^{n-1} \frac{1}{n} f\left(\frac{x}{n}\right) = \pi^{n-1} \frac{1}{n} \frac{1}{1 + \frac{x^2}{n^2}} = \pi^{n-1} \frac{n}{n^2 + x^2}.$$

3. a) By definition

$$a_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \cos nx dx, \quad n \geq 0, \quad b_n = \frac{1}{\pi} \int_0^{2\pi} f(x) \sin nx dx, \quad n \geq 1$$

and

$$c_n = \frac{1}{2\pi} \int_0^{2\pi} f(x) e^{-inx} dx, \quad n \in \mathbf{Z}.$$

Hence

$$c_n = \frac{1}{2\pi} \int_0^{2\pi} f(x)e^{-inx} du = \frac{1}{2\pi} \int_0^{2\pi} f(x)(\cos nx - i \sin nx) dx = \frac{1}{2}(a_n - ib_n), \quad n \geq 1$$

and

$$c_{-n} = \frac{1}{2\pi} \int_0^{2\pi} f(x)e^{inx} du = \frac{1}{2\pi} \int_0^{2\pi} f(x)(\cos nx + i \sin nx) dx = \frac{1}{2}(a_n + ib_n), \quad n \geq 1.$$

Also  $c_0 = \frac{1}{2}a_0$ . Consequently

$$\begin{aligned} \sum_{n=-N}^N c_n e^{inx} &= c_0 + \sum_{n=1}^N (c_n e^{inx} + c_{-n} e^{-inx}) = \\ &= \frac{1}{2}a_0 + \sum_{n=1}^N \left( \frac{1}{2}(a_n - ib_n)(\cos nx + i \sin nx) + \frac{1}{2}(a_n + ib_n)(\cos nx - i \sin nx) \right) = \\ &= \frac{1}{2}a_0 + \sum_{n=1}^N (a_n \cos nx + b_n \sin nx) \quad \text{for all integers } N \geq 1 \text{ and all } x \in \mathbf{R}. \end{aligned}$$

b) The complex Fourier series coefficients of  $f(x)$  are

$$\begin{aligned} c_n &= \frac{1}{2\pi} \int_0^{2\pi} f(x)e^{-inx} dx = \frac{1}{2\pi} \int_0^{2\pi} e^{-iax} e^{-inx} dx = \frac{1}{2\pi} \int_0^{2\pi} e^{-i(n+a)x} dx = \\ &= \frac{1}{2\pi} \left[ \frac{e^{-i(n+a)x}}{-i(n+a)} \right]_0^{2\pi} = \frac{1}{2\pi i} (1 - e^{-i2\pi a}) \frac{1}{n+a}, \quad n \in \mathbf{Z}. \end{aligned}$$

Hence

$$f(x) \sim \sum_{n=-\infty}^{\infty} \frac{1}{2\pi i} (1 - e^{-i2\pi a}) \frac{1}{n+a} e^{inx}$$

is the complex Fourier series expansion of  $f(x)$ . Let  $a_n$ ,  $n \geq 0$ , and  $b_n$ ,  $n \geq 1$ , be the Fourier series coefficients of  $f(x)$ . The function  $f(x)$  is piecewise  $C^1$  on  $\mathbf{R}$ . The Dirichlet convergence theorem for Fourier series therefore shows that the Fourier series of  $f(x)$  converges for all  $x \in \mathbf{R}$  and that its sum is equal to  $\frac{1}{2}(f(x+) + f(x-))$  for all  $x \in \mathbf{R}$ . That is  $\frac{1}{2}a_0 + \sum_{n=1}^{\infty} (a_n \cos nx + b_n \sin nx) = \frac{1}{2}(f(x+) + f(x-))$  for all  $x \in \mathbf{R}$ , or equivalently

$$\frac{1}{2}a_0 + \sum_{n=1}^N (a_n \cos nx + b_n \sin nx) \rightarrow \frac{1}{2}(f(x+) + f(x-)) \quad \text{as } N \rightarrow \infty \text{ for all } x \in \mathbf{R}.$$

Using a) we get

$$(4) \quad \sum_{n=-N}^N \frac{1}{2\pi i} (1 - e^{-i2\pi a}) \frac{1}{n+a} e^{inx} \rightarrow \frac{1}{2}(f(x+) + f(x-)) \quad \text{as } N \rightarrow \infty \text{ for all } x \in \mathbf{R}.$$

Applying (4) with  $x = 0$  we get

$$\sum_{n=-N}^N \frac{1}{2\pi i} (1 - e^{-i2\pi a}) \frac{1}{n+a} \rightarrow \frac{1}{2} (1 + e^{-i2\pi a}) \quad \text{as } N \rightarrow \infty.$$

But

$$\pi i \frac{1 + e^{-i2\pi a}}{1 - e^{-i2\pi a}} = \pi i \frac{e^{i\pi a} + e^{-i\pi a}}{e^{i\pi a} - e^{-i\pi a}} = \pi i \frac{2 \cos \pi a}{2i \sin \pi a} = \frac{\pi}{\cot \pi a}.$$

Hence

$$\sum_{n=-N}^N \frac{1}{n+a} \rightarrow \frac{\pi}{\cot \pi a} \quad \text{as } N \rightarrow \infty.$$

4. The corresponding Sturm-Liouville eigenvalue problem is  $X''(x) + \lambda X(x) = 0$ ,  $0 < x < 1$ , with boundary values  $X(0) = 0$  and  $X'(1) = 0$ . The theory of Sturm-Liouville eigenvalue problems shows that any eigenvalue to this eigenvalue problem must be  $\geq 0$ . The value  $\lambda = 0$  gives  $X(x) = 0$  so  $\lambda = 0$  is not an eigenvalue. Now suppose  $\lambda > 0$ . Then  $\lambda = \alpha^2$  where  $\alpha > 0$ , and  $X''(x) + \lambda X(x) = 0$  gives  $X(x) = A \sin \alpha x + B \cos \alpha x$ . The boundary value  $X(0) = 0$  implies  $B = 0$ . The boundary value  $X'(1) = 0$  then shows that there is a non-zero solution  $X(x)$  if and only if  $\cos \alpha = 0$ . Since  $\alpha > 0$  this happens if and only if  $\alpha = (n + \frac{1}{2})\pi$  and  $n = 0, 1, \dots$ . The Sturm-Liouville problem here therefore has the eigenvalues  $\lambda_n = (n + \frac{1}{2})^2 \pi^2$ ,  $n = 0, 1, \dots$  and the corresponding eigenfunctions  $X_n(x) = \sin(n + \frac{1}{2})\pi x$ ,  $n = 0, 1, \dots$ . The eigenfunctions constitute an orthogonal system with respect to the inner product  $\langle f, g \rangle = \int_0^1 f(x)\overline{g(x)} dx$ . The corresponding norm is  $\|f\| = \sqrt{\langle f, f \rangle}$ . Since  $\|X_n\|^2 = \int_0^1 \sin^2(n + \frac{1}{2})\pi x dx = \frac{1}{2}$  for  $n = 0, 1, \dots$  the Fourier series expansion of a function  $f(x)$  on  $0 \leq x \leq 1$  with respect to the orthogonal system  $X_n(x)$   $n = 0, 1, 2, \dots$  is

$$f(x) \sim \sum_{n=0}^{\infty} \alpha_n X_n(x) \quad \text{where} \quad \alpha_n = 2 \int_0^1 f(x) X_n(x) dx \quad n = 0, 1, 2, \dots$$

We now expand the sought function  $u(x, y)$  in a Fourier series with respect to the orthogonal system  $X_n(x)$   $n = 0, 1, 2, \dots$ . From the general Sturm-Liouville theory we know that  $u(x, y)$  then is equal to its Fourier series, that is

$$(5) \quad u(x, y) = \sum_{n=0}^{\infty} Y_n(y) X_n(x) \quad \text{where} \quad Y_n(y) = 2 \int_0^1 u(x, y) X_n(x) dx \quad n = 0, 1, 2, \dots$$

We also know that

$$u_{xx}(x, y) \sim - \sum_{n=0}^{\infty} \pi^2 \left(n + \frac{1}{2}\right)^2 Y_n(y) X_n(x) \quad \text{and} \quad u_{yy}(x, y) \sim \sum_{n=0}^{\infty} Y_n''(y) X_n(x)$$

so

$$u_{xx}(x, y) + u_{yy}(x, y) \sim \sum_{n=0}^{\infty} \left( Y_n''(y) - \pi^2 \left(n + \frac{1}{2}\right)^2 Y_n(y) \right) X_n(x).$$

But  $u_{xx}(x, y) + u_{yy}(x, y) = 0$ . Hence

$$(6) \quad Y_n''(y) - \pi^2 \left(n + \frac{1}{2}\right)^2 Y_n(y) = 0, \quad n = 0, 1, \dots$$

From (5) we further get

$$(7) \quad u(x, 0) = \sum_{n=0}^{\infty} Y_n(0) X_n(x) \quad \text{and} \quad u(x, 1) = \sum_{n=0}^{\infty} Y_n(1) X_n(x).$$

We know that

$$(8) \quad u(x, 0) = u(x, 1) = x^2 - 2x.$$

In order to be able to exploit that we expand  $f(x) = x^2 - 2x$  in a Fourier series with respect to the orthogonal system  $X_n(x)$ ,  $n = 0, 1, \dots$ . The Fourier series coefficients in that expansion are

$$\begin{aligned} 2 \int_0^1 f(x) X_n(x) dx &= 2 \int_0^1 (x^2 - 2x) \sin \left( n + \frac{1}{2} \right) \pi x dx = \\ &= 2 \left( \left[ (x^2 - 2x) \left( -\frac{\cos \left( n + \frac{1}{2} \right) \pi x}{\left( n + \frac{1}{2} \right) \pi} \right) \right]_0^1 + \frac{2}{\left( n + \frac{1}{2} \right) \pi} \int_0^1 (x - 1) \cos \left( n + \frac{1}{2} \right) \pi x dx \right) = \\ &= \frac{4}{\left( n + \frac{1}{2} \right) \pi} \left( \left[ (x - 1) \frac{\sin \left( n + \frac{1}{2} \right) \pi x}{\left( n + \frac{1}{2} \right) \pi} \right]_0^1 - \frac{1}{\left( n + \frac{1}{2} \right) \pi} \int_0^1 \sin \left( n + \frac{1}{2} \right) \pi x dx \right) = \\ &= -\frac{4}{\left( n + \frac{1}{2} \right)^2 \pi^2} \left[ -\frac{\cos \left( n + \frac{1}{2} \right) \pi x}{\left( n + \frac{1}{2} \right) \pi} \right]_0^1 = \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3}, \quad n = 0, 1, \dots \end{aligned}$$

Thus

$$(9) \quad x^2 - 2x \sim \sum_{n=0}^{\infty} \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3} X_n(x).$$

Comparison of (7), (8) and (9) gives

$$(10) \quad Y_n(0) = Y_n(1) = \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3}, \quad n = 0, 1, \dots$$

From (6) we get

$$Y_n(y) = A_n e^{(n+\frac{1}{2})\pi y} + B_n e^{-(n+\frac{1}{2})\pi y} \quad n = 0, 1, \dots$$

Using (10) we then get

$$A_n + B_n = \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3} \quad \text{and} \quad A_n e^{(n+\frac{1}{2})\pi} + B_n e^{-(n+\frac{1}{2})\pi} = \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3} \quad n = 0, 1, \dots$$

which gives

$$A_n = \frac{1}{1 + e^{(n+\frac{1}{2})\pi}} \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3} \quad \text{and} \quad B_n = \frac{1}{1 + e^{-(n+\frac{1}{2})\pi}} \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3} \quad n = 0, 1, \dots$$

The sought solution  $u(x, y) = \sum_{n=0}^{\infty} Y_n(y) X_n(x)$  therefore is

$$u(x, y) = \sum_{n=0}^{\infty} \frac{4}{\left( n + \frac{1}{2} \right)^3 \pi^3} \left( \frac{e^{(n+\frac{1}{2})\pi y}}{1 + e^{(n+\frac{1}{2})\pi}} + \frac{e^{-(n+\frac{1}{2})\pi y}}{1 + e^{-(n+\frac{1}{2})\pi}} \right) \sin \left( n + \frac{1}{2} \right) \pi x.$$

5. For  $u \neq 0$  we get

$$\begin{aligned} \frac{3}{4} \int_0^1 \cos ux \, dx - \frac{1}{4} \int_0^3 \cos ux \, dx &= \frac{3}{4} \left[ \frac{\sin ux}{u} \right]_{x=0}^{x=1} - \frac{1}{4} \left[ \frac{\sin ux}{u} \right]_{x=0}^{x=3} = \\ &= \frac{1}{u} \left( \frac{3}{4} \sin u - \frac{1}{4} \sin 3u \right). \end{aligned}$$

But

$$\begin{aligned} \sin^3 u &= \left( \frac{e^{iu} - e^{-iu}}{2i} \right)^3 = -\frac{1}{8i} (e^{i3u} - 3e^{iu} + 3e^{-iu} - e^{-i3u}) = \\ &= \frac{3}{4} \frac{e^{iu} - e^{-iu}}{2i} - \frac{1}{4} \frac{e^{i3u} - e^{-i3u}}{2i} = \frac{3}{4} \sin u - \frac{1}{4} \sin 3u \quad \text{for all } u \in \mathbf{R}. \end{aligned}$$

Hence

$$\frac{3}{4} \int_0^1 \cos ux \, dx - \frac{1}{4} \int_0^3 \cos ux \, dx = \frac{\sin^3 u}{u} \quad \text{for } u \neq 0.$$

For  $u = 0$  we get

$$\frac{3}{4} \int_0^1 \cos ux \, dx - \frac{1}{4} \int_0^3 \cos ux \, dx = \frac{3}{4} \int_0^1 dx - \frac{1}{4} \int_0^3 dx = 0.$$

Set

$$f(x) = \begin{cases} 0, & x < -3 \\ -\frac{1}{4}, & -3 \leq x < -1 \\ \frac{1}{2}, & -1 \leq x \leq 1 \\ -\frac{1}{4}, & 1 < x \leq 3 \\ 0, & x > 3. \end{cases}$$

Then  $f(x)$  is an even function and

$$\begin{aligned} \frac{3}{4} \int_0^1 \cos ux \, dx - \frac{1}{4} \int_0^3 \cos ux \, dx &= \frac{1}{2} \int_0^1 \cos ux \, dx - \frac{1}{4} \int_1^3 \cos ux \, dx = \int_0^3 f(x) \cos ux \, dx = \\ &= \frac{1}{2} \int_{-3}^3 f(x) \cos ux \, dx = \frac{1}{2} \int_{-3}^3 f(x) (\cos ux - i \sin ux) \, dx = \frac{1}{2} \int_{-3}^3 f(x) e^{-iux} \, dx = \\ &= \pi \frac{1}{2\pi} \int_{-3}^3 f(x) e^{-iux} \, dx = \pi \frac{1}{2\pi} \int_{-\infty}^{\infty} f(x) e^{-iux} \, dx = \pi \hat{f}(u). \end{aligned}$$

That is

$$(11) \quad \frac{\sin^3 u}{u} = \pi \hat{f}(u) \quad \text{if } u \neq 0.$$

For an arbitrary number  $T > 0$ , an arbitrary  $x \in \mathbf{R}$  and with the use of (11) we get

$$(12) \quad \int_0^T \frac{\sin^3 u}{u} \cos ux \, dx = \frac{1}{2} \int_{-T}^T \frac{\sin^3 u}{u} \cos ux \, dx = \frac{1}{2} \int_{-T}^T \frac{\sin^3 u}{u} (\cos ux + i \sin ux) \, dx =$$

$$= \frac{1}{2} \int_{-T}^T \frac{\sin^3 u}{u} e^{iux} dx = \frac{\pi}{2} \int_{-T}^T \hat{f}(u) e^{iux} dx.$$

But  $f(x)$  is piecewise  $C^1$  on  $\mathbf{R}$  and therefore

$$(13) \quad \int_{-T}^T \hat{f}(u) e^{iux} dx \rightarrow \frac{1}{2}(f(x+) + f(x-)) \quad \text{as } T \rightarrow \infty \text{ for each } x \in \mathbf{R}$$

by the inversion formula for the Fourier transform. Combining (12) and (13) it follows that

$$\int_0^T \frac{\sin^3 u}{u} \cos ux dx \rightarrow \frac{\pi}{4}(f(x+) + f(x-)) \quad \text{as } T \rightarrow \infty \text{ for each } x \in \mathbf{R}.$$

That is the generalized integral

$$\int_0^\infty \frac{\sin^3 u}{u} \cos ux dx$$

is convergent for each  $x \in \mathbf{R}$  and

$$\int_0^\infty \frac{\sin^3 u}{u} \cos ux dx = \frac{\pi}{4}(f(x+) + f(x-)) = \begin{cases} 0, & x < -3 \\ -\frac{\pi}{16}, & x = -3 \\ -\frac{\pi}{8}, & -3 < x < -1 \\ \frac{\pi}{16}, & x = -1 \\ \frac{\pi}{8}, & -1 < x < 1 \\ \frac{\pi}{16}, & x = 1 \\ -\frac{\pi}{8}, & 1 < x < 3 \\ -\frac{\pi}{16}, & x = 3 \\ 0, & x > 3. \end{cases}$$

Finally using (11) and the Plancherel identity we get

$$\begin{aligned} \int_0^\infty \frac{\sin^6 u}{u^2} du &= \frac{1}{2} \int_{-\infty}^\infty \frac{\sin^6 u}{u^2} du = \frac{1}{2} \int_{-\infty}^\infty \left| \frac{\sin^3 u}{u} \right|^2 du = \frac{\pi^2}{2} \int_{-\infty}^\infty |\hat{f}(u)|^2 du = \\ &= \frac{\pi}{4} \int_{-\infty}^\infty |f(x)|^2 dx = \frac{\pi}{4} \int_{-3}^3 |f(x)|^2 dx = \frac{\pi}{2} \int_0^3 |f(x)|^2 dx = \\ &= \frac{\pi}{2} \left( \frac{1}{4} + 2 \cdot \frac{1}{16} \right) = \frac{3\pi}{16}. \end{aligned}$$



6. a) Let  $f(x) \in A$  and let  $\sum_{n=-\infty}^{\infty} c_n e^{inx}$  be the complex Fourier series expansion of  $f(x)$ . Then  $\sum_{n=-\infty}^{\infty} |c_n| < \infty$  by the definition of  $A$ . But  $|c_n e^{inx}| = |c_n|$  for all  $n \in \mathbf{Z}$  and all  $x \in \mathbf{R}$ . Hence  $\sum_{n=-\infty}^{\infty} c_n e^{inx}$  converges absolutely and uniformly on  $\mathbf{R}$ . Set  $g(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx}$  for  $x \in \mathbf{R}$ . Each function  $c_n e^{inx}$  is continuous on  $\mathbf{R}$ . An infinite sum of continuous functions on  $\mathbf{R}$  converging uniformly on  $\mathbf{R}$  is continuous on  $\mathbf{R}$ . Hence  $g(x)$  is continuous on  $\mathbf{R}$ . The function  $g(x)$  is also  $2\pi$ -periodic on  $\mathbf{R}$  since each function  $c_n e^{inx}$  is  $2\pi$ -periodic on  $\mathbf{R}$ . The  $m$ -th complex Fourier series coefficient of  $g(x)$  is

$$\begin{aligned} \frac{1}{2\pi} \int_0^{2\pi} g(x) e^{-imx} dx &= \frac{1}{2\pi} \int_0^{2\pi} \left( \sum_{n=-\infty}^{\infty} c_n e^{inx} \right) e^{-imx} dx = \\ &= \sum_{n=-\infty}^{\infty} c_n \frac{1}{2\pi} \int_0^{2\pi} e^{i(n-m)x} dx = c_m. \end{aligned}$$

The interchange of integration and summation above is justified by the uniform convergence of the sum on  $\mathbf{R}$ . Above we have also used that

$$\frac{1}{2\pi} \int_0^{2\pi} e^{i(n-m)x} dx = \begin{cases} 1 & \text{if } n = m \\ 0 & \text{if } n \neq m. \end{cases}$$

The functions  $f(x)$  and  $g(x)$  thus have the same complex Fourier series coefficients. Consequently all complex Fourier series coefficients of  $f(x) - g(x)$  are zero. Hence

$$\int_0^{2\pi} |f(x) - g(x)|^2 dx = 0$$

by the Parseval theorem. Since  $f(x) - g(x)$  is continuous on  $\mathbf{R}$  ( $f(x)$  is continuous on  $\mathbf{R}$  since  $f(x) \in A$ ) it follows that  $f(x) = g(x)$  on  $0 \leq x \leq 2\pi$  and consequently on  $\mathbf{R}$  because of the  $2\pi$ -periodicity of both functions. Hence  $f(x) = g(x) = \sum_{n=-\infty}^{\infty} c_n e^{inx}$  for all  $x \in \mathbf{R}$ .

b) Let  $f(x)$  be a complex-valued  $2\pi$ -periodic continuous function  $f(x)$  on  $\mathbf{R}$  and assume that  $f(x)$  has a piecewise continuous derivative on  $\mathbf{R}$ . Let  $f(x) \sim \sum_{n=-\infty}^{\infty} c_n e^{inx}$  be the complex Fourier series expansion of  $f(x)$ . The  $n$ -th complex Fourier series coefficient of  $f'(x)$  then is

$$\begin{aligned} \frac{1}{2\pi} \int_0^{2\pi} f'(x) e^{-inx} dx &= \frac{1}{2\pi} \left( [f(x) e^{-inx}]_0^{2\pi} - \int_0^{2\pi} f(x) (-in e^{-inx}) dx \right) = \\ &= in \frac{1}{2\pi} \int_0^{2\pi} f(x) e^{-inx} dx = in c_n. \end{aligned}$$

Hence  $f'(x) \sim \sum_{n=-\infty}^{\infty} in c_n e^{inx}$  is the complex Fourier series expansion of  $f'(x)$ . Using the Cauchy-Schwarz inequality and Bessel's inequality for  $f'(x)$  we get

$$\begin{aligned} \sum_{\substack{-N \leq n \leq N \\ n \neq 0}} |c_n| &= \sum_{\substack{-N \leq n \leq N \\ n \neq 0}} \frac{1}{|n|} |n c_n| \leq \left( \sum_{\substack{-N \leq n \leq N \\ n \neq 0}} \frac{1}{n^2} \right)^{1/2} \left( \sum_{\substack{-N \leq n \leq N \\ n \neq 0}} n^2 |c_n|^2 \right)^{1/2} \leq \\ &\leq \left( 2 \sum_{n=1}^{\infty} \frac{1}{n^2} \right)^{1/2} \left( \frac{1}{2\pi} \int_0^{2\pi} |f'(x)|^2 dx \right)^{1/2} \quad \text{for all integers } N \geq 1. \end{aligned}$$

Thus  $\sum_{n=-\infty}^{\infty} |c_n| < \infty$  since  $\sum_{n=1}^{\infty} \frac{1}{n^2} < \infty$ . Consequently  $f(x) \in A$ .

c) We first note that

$$h(x + 2\pi) = \int_{-\pi}^{\pi} f(x + 2\pi - y)g(y) dy = \int_{-\pi}^{\pi} f(x - y)g(y) dy = h(x) \quad \text{for all } x \in \mathbf{R}$$

since  $f(x)$  is  $2\pi$ -periodic. Hence  $h(x)$  is  $2\pi$ -periodic. We now show that  $h(x)$  is continuous on  $\mathbf{R}$ . This can be proved in different ways. Here we give a Fourier analysis proof. Fix an arbitrary  $x \in \mathbf{R}$  and let  $\delta \in \mathbf{R}$ . Then

$$\begin{aligned} (14) \quad |h(x + \delta) - h(x)| &= \left| \int_{-\pi}^{\pi} f(x + \delta - y)g(y) dy - \int_{-\pi}^{\pi} f(x - y)g(y) dy \right| = \\ &= \left| \int_{-\pi}^{\pi} (f(x - y + \delta) - f(x - y))g(y) dy \right| \leq \int_{-\pi}^{\pi} |f(x - y + \delta) - f(x - y)| |g(y)| dy \leq \\ &\leq \left( \int_{-\pi}^{\pi} |f(x - y + \delta) - f(x - y)|^2 dy \right)^{1/2} \left( \int_{-\pi}^{\pi} |g(y)|^2 dy \right)^{1/2} \end{aligned}$$

using the triangle inequality and the Cauchy-Schwarz inequality for integrals. Also

$$(15) \quad \int_{-\pi}^{\pi} |f(x - y + \delta) - f(x - y)|^2 dy = \int_{x-\pi}^{x+\pi} |f(u + \delta) - f(u)|^2 du = \int_{-\pi}^{\pi} |f(u + \delta) - f(u)|^2 du$$

since  $f(x)$  is  $2\pi$ -periodic. Let  $f(u) \sim \sum_{n=-\infty}^{\infty} a_n e^{inu}$  be the complex Fourier series expansion of  $f(u)$ . Then  $f(u + \delta) - f(u) \sim \sum_{n=-\infty}^{\infty} a_n (e^{in\delta} - 1) e^{inu}$  is the complex Fourier series expansion of  $f(u + \delta) - f(u)$ . Applying the Parseval identity to  $f(u)$  and  $f(u + \delta) - f(u)$  we get

$$(16) \quad \frac{1}{2\pi} \int_{-\pi}^{\pi} |f(u)|^2 du = \sum_{n=-\infty}^{\infty} |a_n|^2$$

and

$$(17) \quad \frac{1}{2\pi} \int_{-\pi}^{\pi} |f(u + \delta) - f(u)|^2 du = \sum_{n=-\infty}^{\infty} |a_n|^2 |e^{in\delta} - 1|^2.$$

Since  $|e^{in\delta} - 1| \leq 2$  for all  $n \in \mathbf{Z}$  and all  $\delta \in \mathbf{R}$  it follows from (16) that  $\sum_{n=-\infty}^{\infty} |a_n|^2 |e^{in\delta} - 1|^2$  converges uniformly on  $\mathbf{R}$ . Since each function  $|a_n|^2 |e^{in\delta} - 1|^2$  is a continuous function of  $\delta$  on  $\mathbf{R}$  it follows that  $\sum_{n=-\infty}^{\infty} |a_n|^2 |e^{in\delta} - 1|^2$  is a continuous function of  $\delta$  on  $\mathbf{R}$ . Hence

$$(18) \quad \sum_{n=-\infty}^{\infty} |a_n|^2 |e^{in\delta} - 1|^2 \rightarrow \sum_{n=-\infty}^{\infty} |a_n|^2 |e^0 - 1|^2 = 0 \quad \text{as } \delta \rightarrow 0.$$

From (14), (15), (17) and (18) we see that  $h(x + \delta) \rightarrow h(x)$  as  $\delta \rightarrow 0$  so  $h(x)$  is continuous on  $\mathbf{R}$ .

We now complete the proof that  $h(x) \in A$  by proving the absolute convergence of the complex Fourier series coefficients of  $h(x)$ . Let  $c_n$ ,  $n \in \mathbf{Z}$ , be the complex Fourier series coefficients of

$h(x)$ , let  $b_n$ ,  $n \in \mathbf{Z}$ , be the complex Fourier series coefficients of  $g(x)$ , and as above let  $a_n$ ,  $n \in \mathbf{Z}$ , be the complex Fourier series coefficients of  $f(x)$ . Then

$$\begin{aligned}
c_n &= \frac{1}{2\pi} \int_{-\pi}^{\pi} h(x) e^{-inx} dx = \frac{1}{2\pi} \int_{-\pi}^{\pi} \left( \int_{-\pi}^{\pi} f(x-y) g(y) dy \right) e^{-inx} dx = \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \left( \int_{-\pi}^{\pi} f(x-y) e^{-in(x-y)} dx \right) g(y) e^{-iny} dy = \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \left( \int_{-\pi+y}^{\pi+y} f(x) e^{-in(x)} dx \right) g(y) e^{-iny} dy = \\
&= \frac{1}{2\pi} \int_{-\pi}^{\pi} \left( \int_{-\pi}^{\pi} f(x) e^{-in(x)} dx \right) g(y) e^{-iny} dy = \\
&= 2\pi \left( \frac{1}{2\pi} \int_{-\pi}^{\pi} f(x) e^{-in(x)} dx \right) \left( \frac{1}{2\pi} \int_{-\pi}^{\pi} g(y) e^{-iny} dy \right) = 2\pi a_n b_n \text{ for all } n \in \mathbf{Z}.
\end{aligned}$$

Using this, the Cauchy-Schwarz inequality and Bessel's inequality for  $f(x)$  and  $g(x)$  we get

$$\begin{aligned}
\sum_{n=-N}^N |c_n| &= 2\pi \sum_{n=-N}^N |a_n| |b_n| \leq 2\pi \left( \sum_{n=-N}^N |a_n|^2 \right)^{1/2} \left( \sum_{n=-N}^N |b_n|^2 \right)^{1/2} \leq \\
&\leq \left( \int_{-\pi}^{\pi} |f(x)|^2 dx \right)^{1/2} \left( \int_{-\pi}^{\pi} |g(x)|^2 dx \right)^{1/2} \text{ for all integers } N \geq 1.
\end{aligned}$$

Hence  $\sum_{n=-\infty}^{\infty} |c_n| < \infty$ . Thus  $h(x) \in A$ .